

Simulating Polymer Systems on GPU

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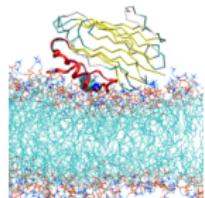


IMPRS School 2012: *GPU Computing – Methods and Applications in the Natural Sciences*

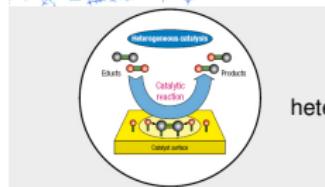
Wroclaw, Poland, 29 October – 02 November 2012

Polymers adsorbing to substrates

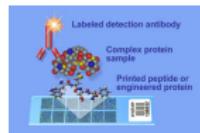
Applications



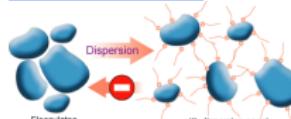
protein-membrane interactions



heterogeneous catalysis



protein microarray (www.arrayit.com)
detection of individual molecules
(hampered by unspecific binding)

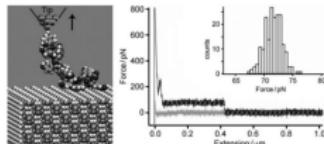


steric colloidal stabilisation
(paints, inks, pharmaceuticals,
food dispersions ...)

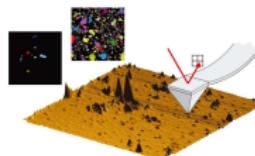


lubrication

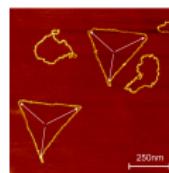
Experiments



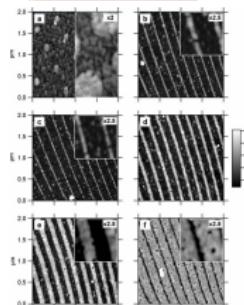
ChemPhysChem **10**, 2795 (2009).
AFM-based single-molecule method,
can determine adhesion free energies



e.g. peptide covered fraction of
substrate [Bachmann/Goede/Beck-
Sickinger/Grundmann/Irbäck/WJ
Angew. Chem. Int. Ed. **49**, 9530 (2010)]



Nano Lett. **4**, 577 (2004).
STM for imaging and manipulating
dsDNA



J. Am. Chem. Soc. **127**, 4323 (2005).
antigen 69k adsorbed on
alkylsilane stripes printed on
a background of oligo(ethylene oxide)

Modeling strategies

- Detailed studies of quantitative properties require microscopic approach based on atomistic models with many fine-tuned parameters
- To get a qualitative overview of the generic behaviour and phase diagrams we use here a coarse-grained approach based on minimalistic models

- Minimalistic (lattice) models:
 - Random Walks – Polymers at theta point (1950–1970)
 - Self-Avoiding Random Walks – Polymers in good solvent (1970–1990)
E. Eisenriegler, K. Kremer, K. Binder, J. Chem. Phys. **77** (1982) 6296
 - Self-Interacting Self-Avoiding Random Walks – Polymers in bad (and good) solvent exhibiting also collapsed states (1990–...)

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Simulation strategies

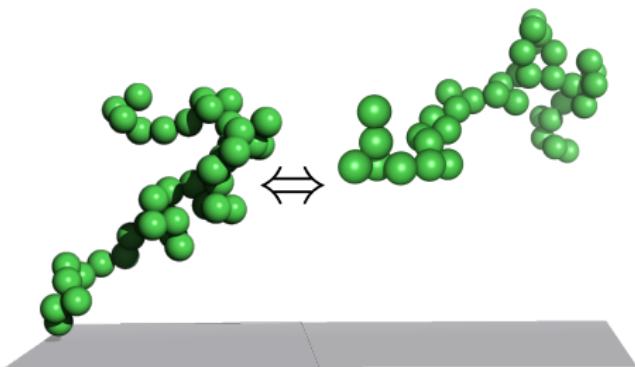
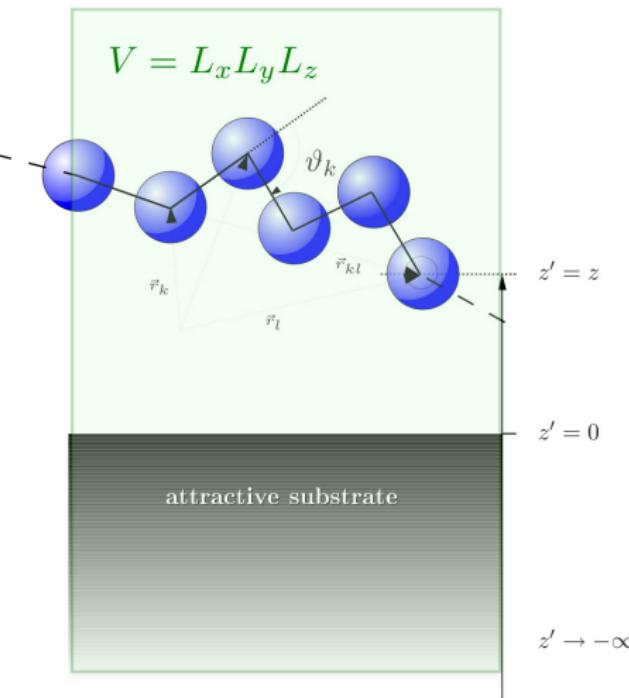
- Off-lattice bead-stick or bead-spring models
- Monte Carlo computer simulations in generalized ensembles:
 - multicanonical
 - parallel tempering
 - ...
- Sophisticated analysis tools:
 - microcanonical quantities
 - finite-size scaling
 - ...

Lattice protein and polymer adsorption (HP model):

M. Bachmann, WJ, Phys. Rev. Lett. **95** (2005) 058102; Phys. Rev. E **73** (2006) 020901(R); Phys. Rev. E **73** (2006) 041802

Flat attractive substrate

model setup

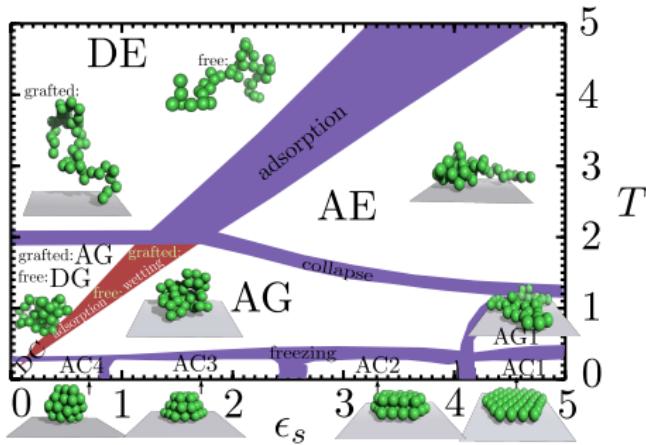


"grafted" versus "free"

bead-stick model

Flat substrate: Phase diagram

substrate attraction strength ϵ_s , temperature T



| pseudophase | typical configuration |
|-------------|-----------------------|
| DE | |
| DG | |
| DC | |
| AE1 | |
| AE2 | |
| AC1 | |
| AG | |
| AC2a | |
| AC2b | |

$$E = 4 \sum_{i=1}^{N-2} \sum_{j=i+2}^N \left(r_{ij}^{-12} - r_{ij}^{-6} \right) + \frac{1}{4} \sum_{i=1}^{N-2} [1 - \cos(\vartheta_i)] + \epsilon_s \sum_{i=1}^N \left(\frac{2}{15} z_i^{-9} - z_i^{-3} \right)$$

M. Möddel, M. Bachmann, WJ, J. Phys. Chem. B **113** (2009) 3314; Phys. Chem. Chem. Phys. **12** (2010) 11548; Macromolecules **44** (2011) 9013

Flat substrate: Method and observables

Method

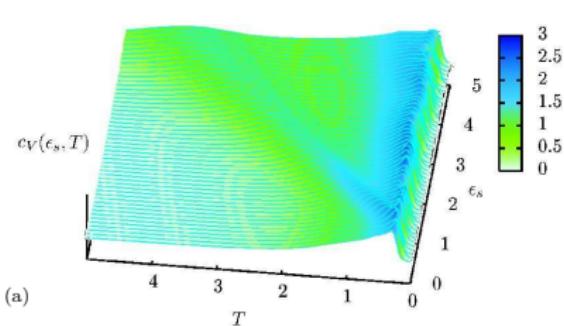
- Multicanonical Monte Carlo computer simulations ($N = 20$ monomers) and parallel tempering ($N = 40$ monomers)

Main observables

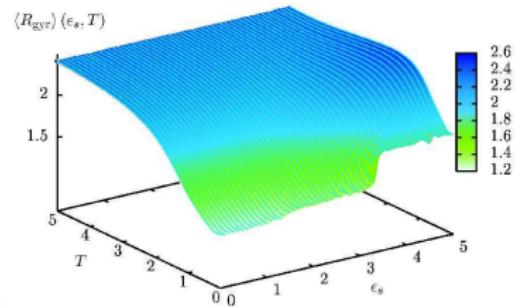
- Energy $\langle E \rangle$ and its individual contributions
- Associated heat capacity $C = \partial \langle E \rangle / \partial T$
- Polymer center-of-mass $\mathbf{r}_{\text{cm}} = N^{-1} \sum_{n=1}^N \mathbf{r}_n$ (average distance to substrate)
- Radius of gyration $\langle R_g \rangle = \left\langle \left[N^{-1} \sum_{n=1}^N (\mathbf{r}_n - \mathbf{r}_{\text{cm}})^2 \right]^{1/2} \right\rangle$ and its components parallel and perpendicular to the substrate
- Polymer-substrate contacts $\langle N_S \rangle$ (counted if $z_i \leq 1.5$)
- Fluctuations $d\langle O \rangle / dT = (\langle OE \rangle - \langle O \rangle \langle E \rangle) / T^2$

Flat substrate: Canonical observables

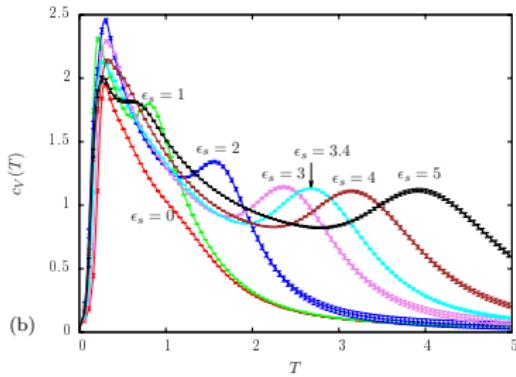
Specific heat and radius of gyration for 20mer



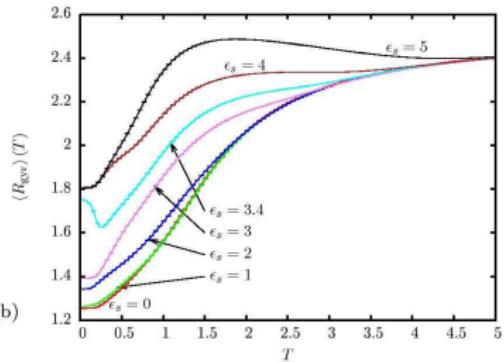
(a)



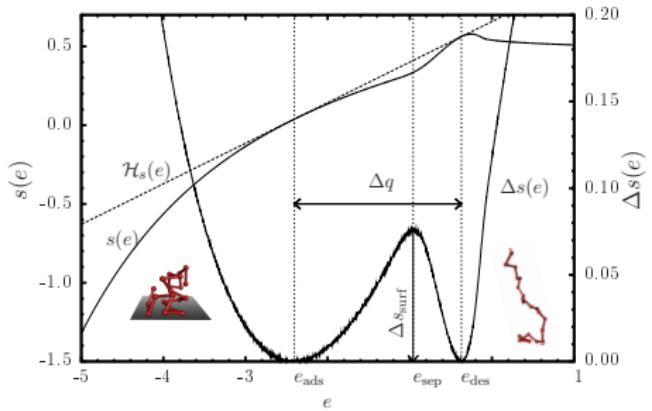
(a)



(b)



Flat substrate: Microcanonical observables



The microcanonical entropy

$$s(E) = k_B \ln g(E)/N$$

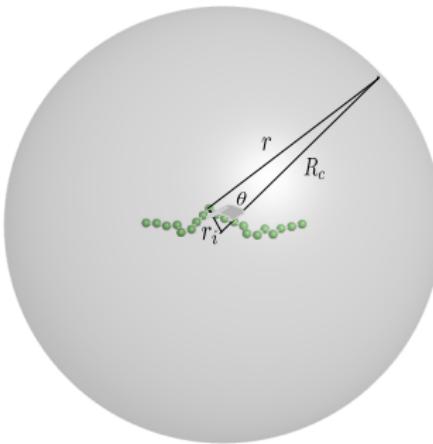
$(Z(\beta) = \sum_E g(E) \exp(-\beta E))$ is a useful signature for first-order like transitions. The Gibbs hull

$$\mathcal{H}_s(e) = s(e_{\text{ads}}) + e(\partial s / \partial e)_{e=e_{\text{ads}}}$$

is the tangent that touches $s(e_{\text{ads}})$ and $s(e_{\text{des}})$. This gives

$$T_{\text{micro}}^{\text{ads}} = \left(\frac{\partial \mathcal{H}_s}{\partial e} \right)^{-1} = \left(\frac{\partial s}{\partial e} \right)^{-1}_{e=e_{\text{ads}}} = \left(\frac{\partial s}{\partial e} \right)^{-1}_{e=e_{\text{des}}}$$

Curved attractive substrate – Spherical cage



Study influence of

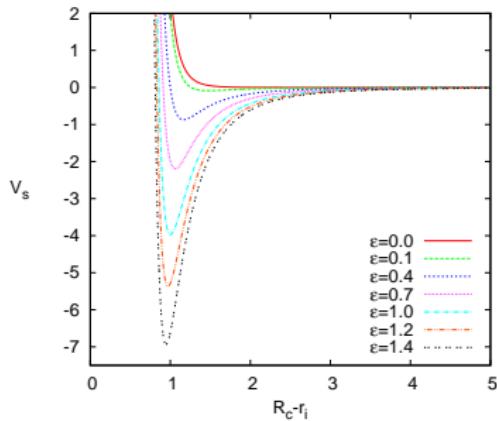
- curved vs flat substrates
- spherical confinement vs half-space

H. Arkın, WJ, Phys. Rev. E **85**, 051802 (2012); J. Phys. Chem. B **116**, 10379 (2012);
Eur. Phys. J. – Special Topics (in print)

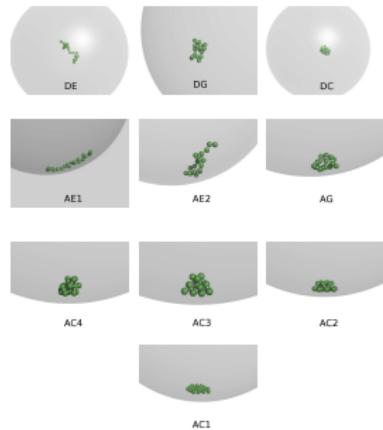
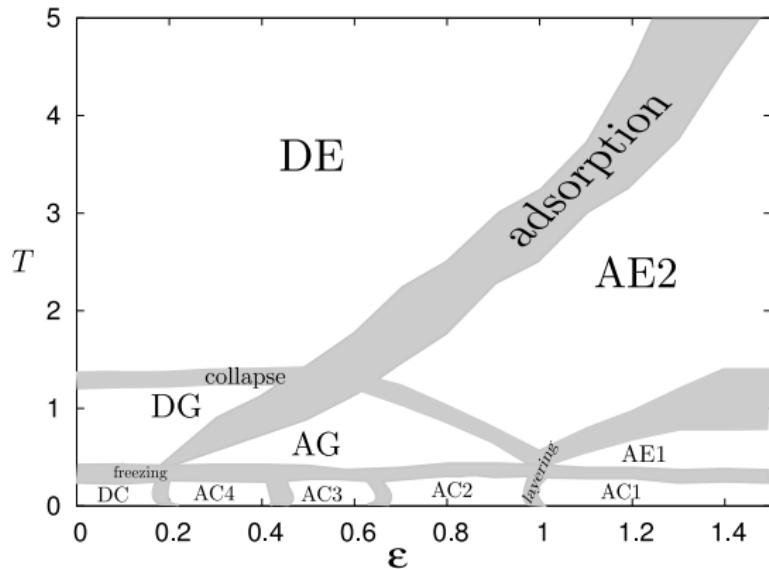
Polymer in spherical cage: Interactions

- Monomer-monomer interaction as before: $12 - 6$ LJ plus (very) weak bending energy
- Monomer-sphere interaction by integrating $12 - 6$ LJ over the surface of the sphere (radius R_c):

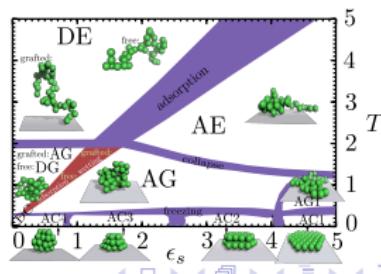
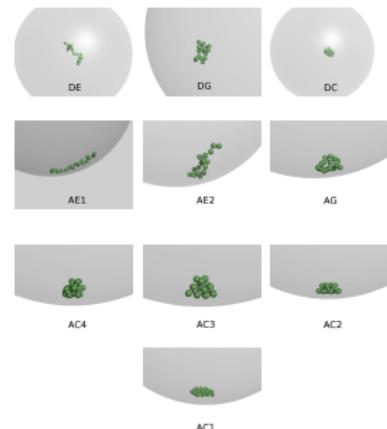
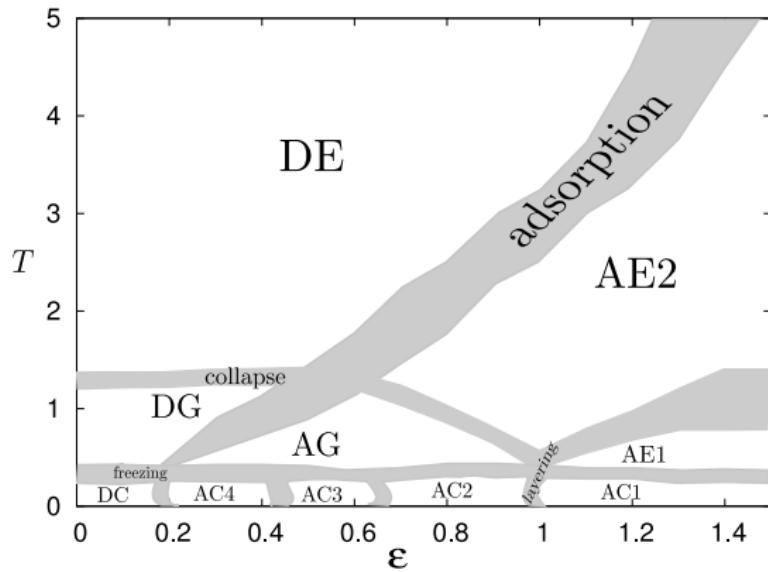
$$V_s = 4 \frac{\pi R_c}{r_i} \left\{ \frac{1}{5} \left[\left(\frac{1}{r_i - R_c} \right)^{10} - \left(\frac{1}{r_i + R_c} \right)^{10} \right] - \frac{\epsilon}{2} \left[\left(\frac{1}{r_i - R_c} \right)^4 - \left(\frac{1}{r_i + R_c} \right)^4 \right] \right\}$$



Polymer in spherical cage: Phase diagram



Polymer in spherical cage: Phase diagram



Similar to phase diagram for flat substrate:

Polymer grafted to fluctuating attractive membrane

Polymer **grafted** to dynamically fluctuating membrane

$$E = E^{\text{mem}} + E^{\text{pol}} + E^{\text{int}}$$

Membrane: Elastic $L_x \times L_y$ mesh with tethering FENE potential between neighboring nodes (plus hard-sphere potential to ensure self-avoidance)

$$V_{\text{FENE}}^{\text{m}}(r) = -\frac{K}{2} R_{\text{m}}^2 \ln \left\{ 1 - [(r - r_0)/R_{\text{m}}]^2 \right\}$$

$$E^{\text{mem}} = \sum_{k=1}^{L_x-1} \sum_{l=1}^{L_y} V_{\text{FENE}}^{\text{m}}(|\mathbf{r}_{k,l} - \mathbf{r}_{k+1,l}|) + \sum_{k=1}^{L_x} \sum_{l=1}^{L_y-1} V_{\text{FENE}}^{\text{m}}(|\mathbf{r}_{k,l} - \mathbf{r}_{k,l+1}|)$$

Polymer: Bead-spring model

$$E^{\text{pol}} = \sum_{i=1}^{N-1} \sum_{j=i+1}^N V_{\text{LJ}}^{\text{pp}}(|\mathbf{r}_i - \mathbf{r}_j|) + \sum_{i=1}^{N-1} V_{\text{FENE}}^{\text{p}}(|\mathbf{r}_i - \mathbf{r}_{i+1}|)$$

$$V_{\text{LJ}}^{\text{pp}}(r) = 4\epsilon_{\text{pp}} \left[(\sigma/r)^{12} - (\sigma/r)^6 \right]; V_{\text{FENE}}^{\text{p}}(r) = -\frac{K}{2} R_{\text{p}}^2 \ln \left\{ 1 - [(r - r_0)/R_{\text{p}}]^2 \right\}$$

Fluctuating membrane: Parameters

Polymer-membrane attraction: Between monomers and mesh nodes

$$V_{\text{LJ}}^{\text{pm}}(r) = 4\epsilon_{\text{pm}} \left[(\sigma/r)^{12} - (\sigma/r)^6 \right]$$

Parameters ($\sigma = r_0/2^{1/6}$; $r_0 = 1$):

Membrane: $L_x = L_y = 27$; $K = 40$; $R_m = 0.1$

Polymer: $N = 13$, grafted at membrane center, $\epsilon_{\text{pp}} = 1$; $K = 40$; $R_p = 0.3$

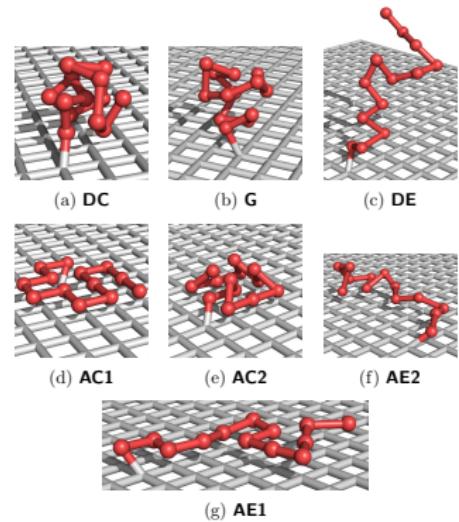
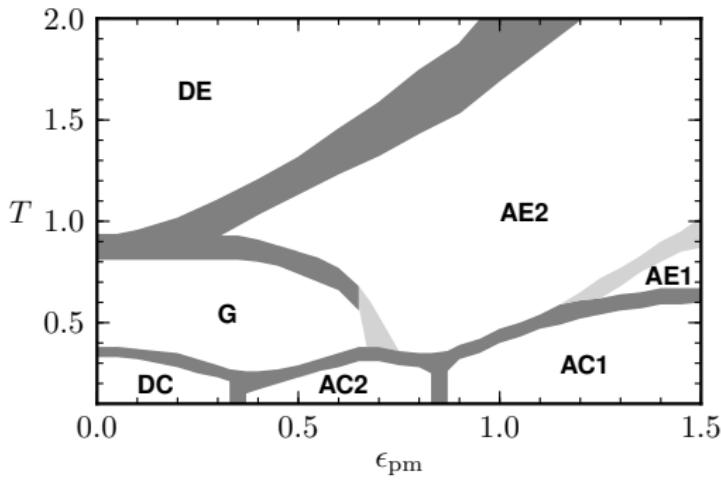
Membrane-polymer attraction parametrized by $\epsilon_{\text{pm}} \in [0, 1.5]$

Stiff membrane (flat, static 2d mesh): Slightly different model than in our recent work on polymer adsorption to a flat substrate, so study this case first

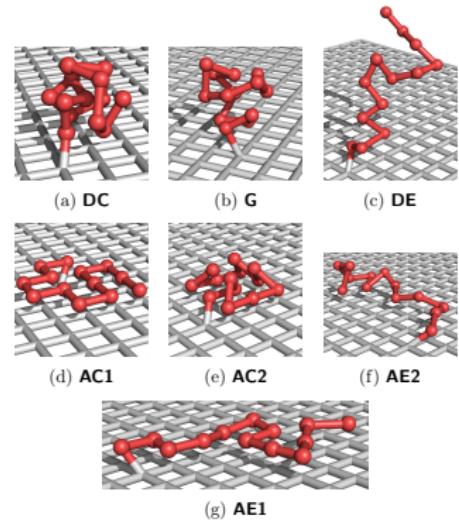
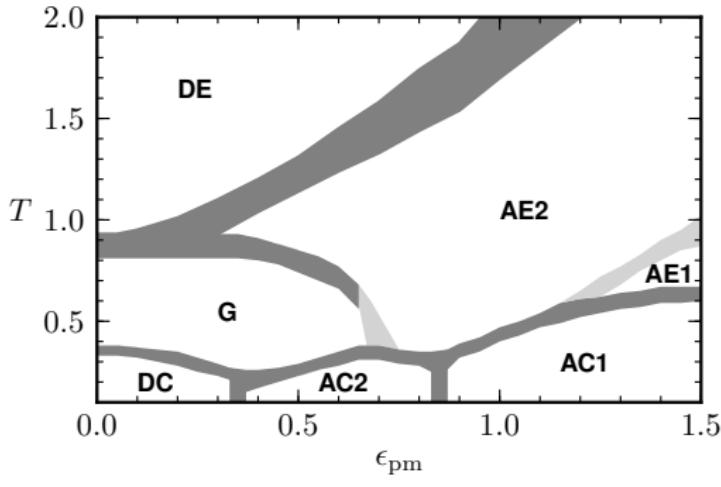
- as reference system for fluctuating membrane
- and for comparison with flat substrate

S. Karalus, WJ, M. Bachmann, Phys. Rev. E **84** (2011) 031803

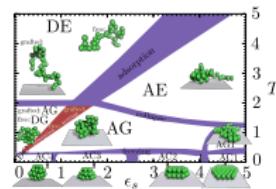
Stiff membrane: Phase diagram



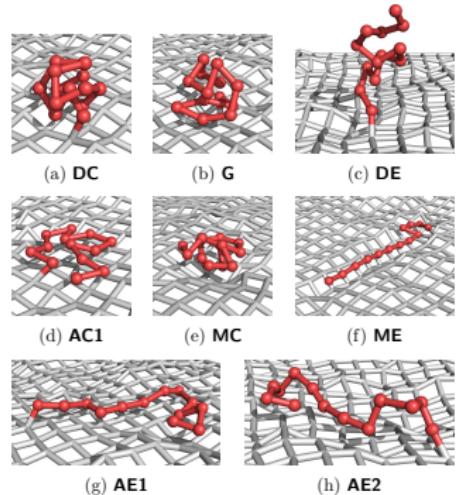
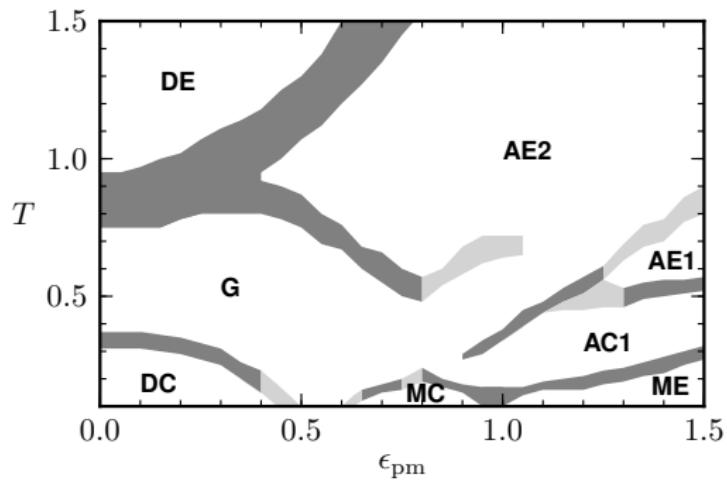
Stiff membrane: Phase diagram



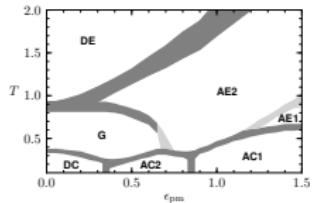
Similar to phase diagram for flat substrate:



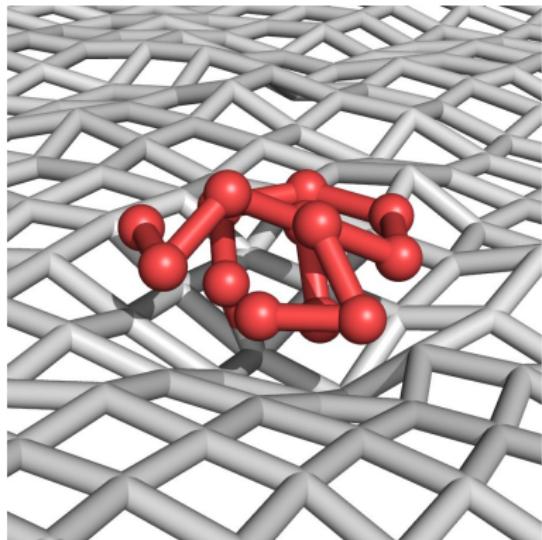
Fluctuating membrane: Phase diagram



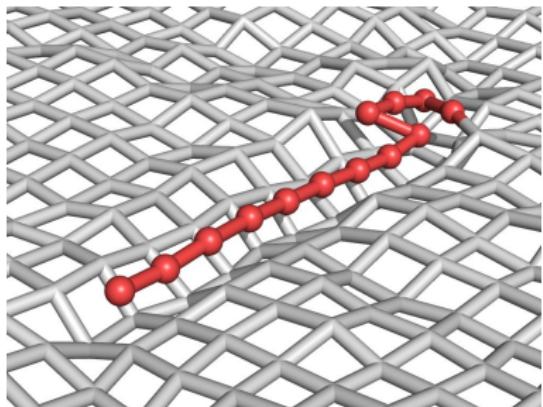
Similar to stiff membrane, but now “embedded” phases:



Fluctuating membrane: “Embedded” phases



Embedded Compact (MC)



Embedded Elongated (ME)

Back-reaction between polymer and membrane fluctuations

Summary and Outlook

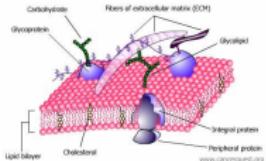
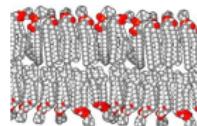
- Adsorption phase diagram of generic polymer model
- Morphologies of low-temperature “frozen” conformations
- Generic phase structure robust and similar to flat substrate system

Summary and Outlook

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Outlook

- Use more realistic polymers (add bending energy for semiflexible (bio-)polymers, simulate all-atom models, ...)
- Improve substrate model (other geometries, patterned surfaces, fluctuating surfaces, better controlled curvature, ...)
- Use atomistic lipid bilayer membranes ...
- ... together with cholesterols, membrane proteins, solvent, and ...

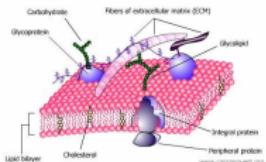
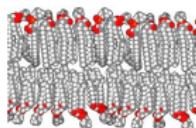


Summary and Outlook

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... something for Wroclaw 2022 !!!

Acknowledgements

Collaborations with

- Michael Bachmann (now Univ. of Georgia, Athens, USA)
- Monika Möddel (now Basycon Consulting, Munich, Germany)
- Handan Arkın-Olgar (on leave from Ankara University, Turkey)
- Steffen Karalus (now University of Cologne, Germany)

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- Alexander von Humboldt Foundation

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Recent Developments in Computational Physics

29 November – 01 December 2012, ITP, Universität Leipzig
www.physik.uni-leipzig.de/~janke/CompPhys12

See you there ?

Bead-Spring Homopolymer Model

Lennard-Jones potential ($\epsilon = 1$, $\sigma = 2^{-1/6}r_0$, $r_0 = 0.7$):

$$E_{\text{LJ}}(r_{ij}) = 4\epsilon \left[\left(\frac{\sigma}{r_{ij}}\right)^{12} - \left(\frac{\sigma}{r_{ij}}\right)^6 \right]$$

Truncated at $r_c = 2.5\sigma$:

$$E_{\text{LJ}}^{\text{mod}}(r_{ij}) = E_{\text{LJ}}(\min(r_{ij}, r_c)) - E_{\text{LJ}}(r_c)$$

Finitely extensible nonlinear elastic (FENE) anharmonic potential ($R = 0.3$, $K = 40$):

$$E_{\text{FENE}}(r_{ii+1}) = -\frac{K}{2}R^2 \log \left[1 - \left(\frac{r_{ii+1} - r_0}{R} \right)^2 \right]$$

Total energy of a conformation $\mathcal{C} = (\mathbf{r}_1, \dots, \mathbf{r}_N)$:

$$E(\mathcal{C}) = \sum_{i < j}^N E_{\text{LJ}}^{\text{mod}}(r_{ij}) + \sum_i^{N-1} E_{\text{FENE}}(r_{ii+1})$$

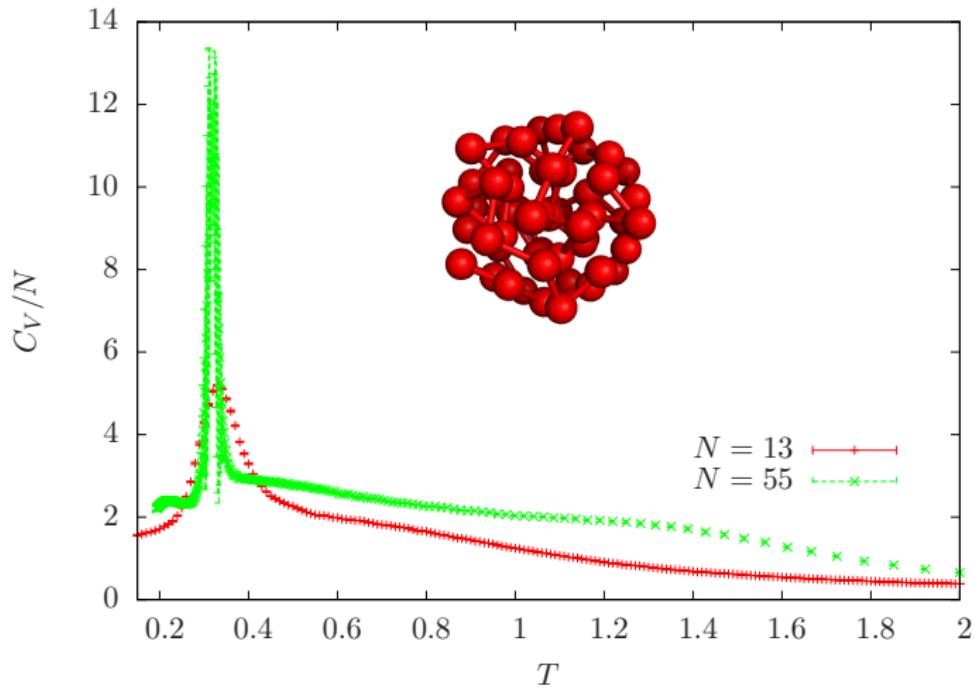
Simulation Method

Parallel tempering (replica exchange) Monte Carlo with n_r replicas:
Exchange conformations of neighboring replicas i and $i + 1$ with probability

$$p = \min(1, \exp[(E_i - E_{i+1})(\beta_i - \beta_{i+1})])$$

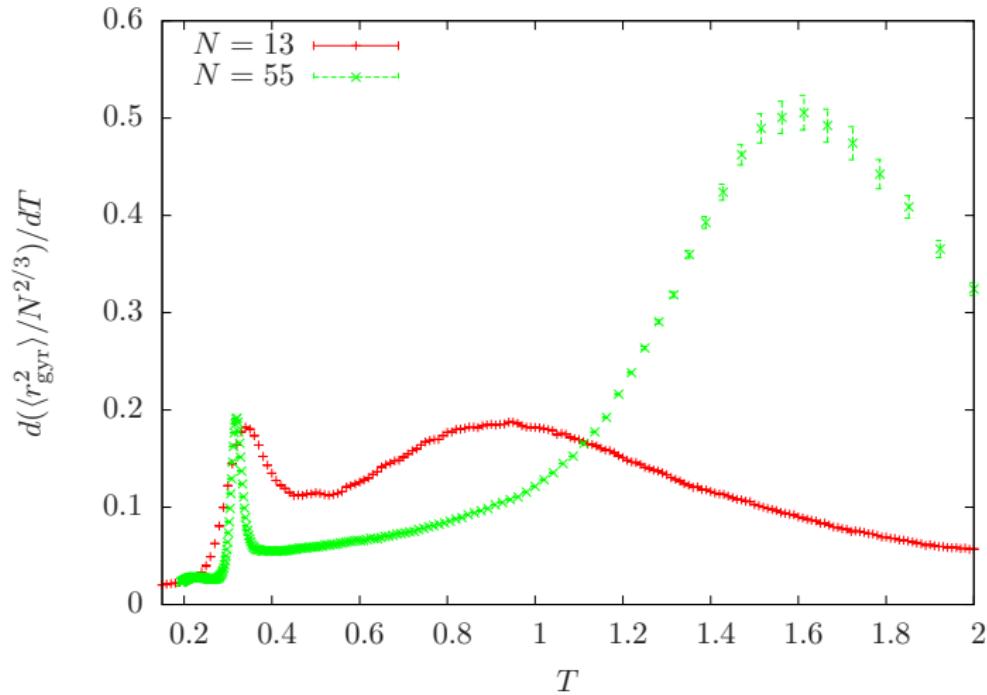
Specific Heat

Specific heat of an elastic polymer chain with lengths $N = 13$ and 55 . The inset shows the icosahedral structure of the 55mer at low T .

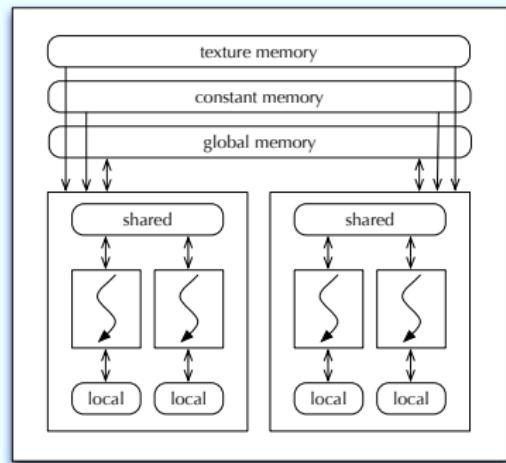
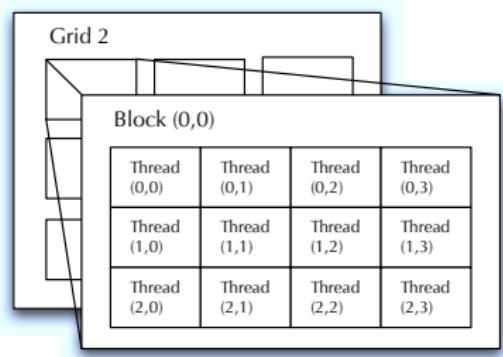
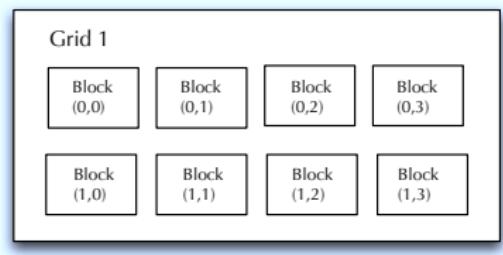


Radius of Gyration

Temperature derivative of mean squared radius of gyration.



GPU: Grid With Thread Blocks and Memory Layout

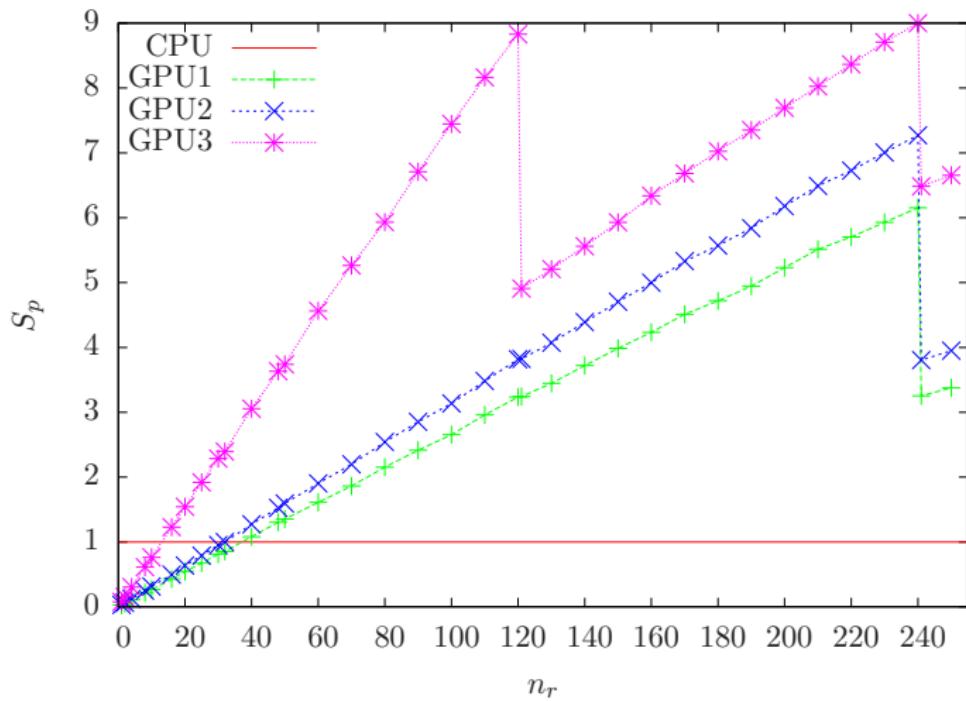


GPU: Specification of used Hardware

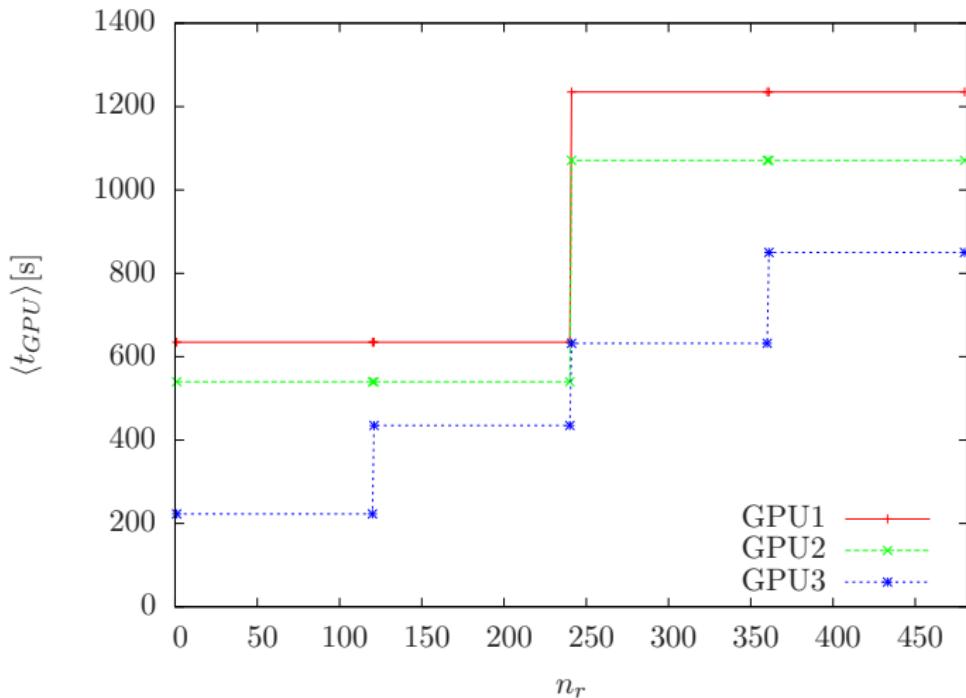
Table: Specifications of the used hardware.

| | reference CPU | GPU1 | GPU2 | GPU3 |
|------------------------|---------------|-------------|---------|--------|
| name | Xeon E5620 | Tesla C1060 | GTX285 | GTX480 |
| # processors | 1 | 30 | 30 | 15 |
| # cores per processor | 4 (1 used) | 8 | 8 | 32 |
| RAM | 16384MB | 4096MB | 1024MB | 1536MB |
| clock speed | 2.4GHz | 1.3GHz | 1.48GHz | 1.4GHz |
| max. threads per block | - | 512 | 512 | 1024 |
| shared memory size | - | 16kB | 16KB | 48kB |
| registers per block | - | 16384 | 16384 | 32768 |

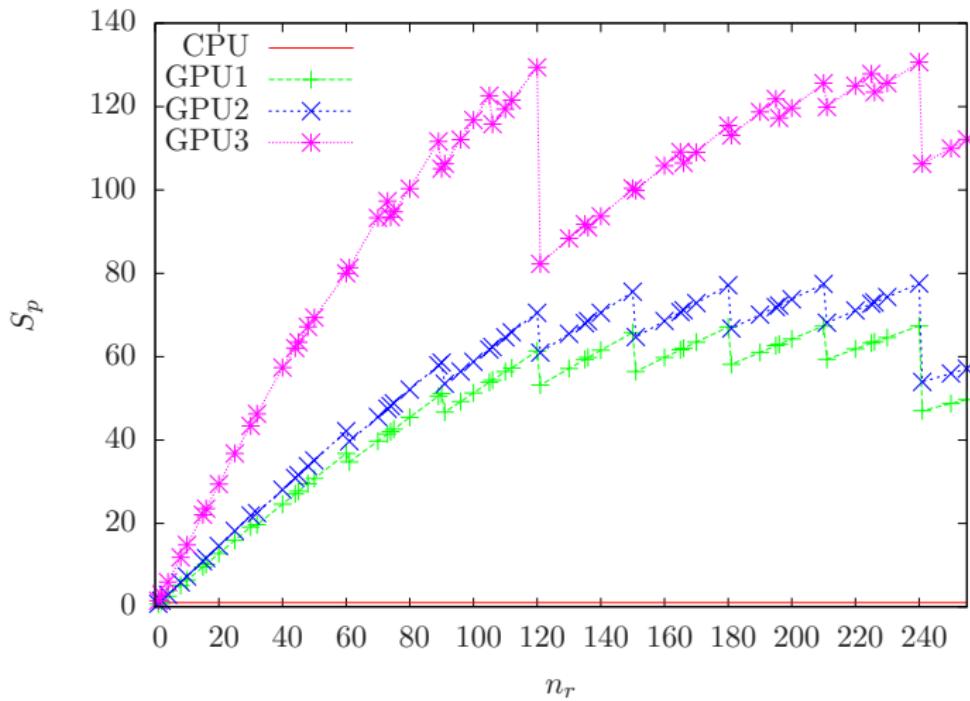
Naive Parallel Tempering Implementation



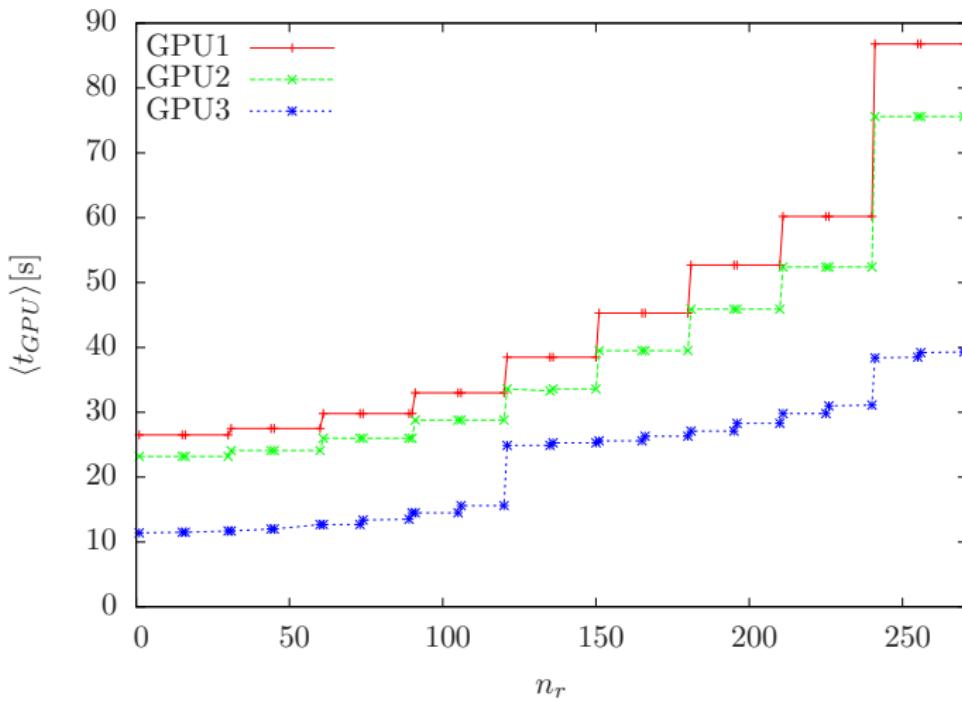
Naive Parallel Tempering Implementation



Improved Parallel Tempering Implementation



Improved Parallel Tempering Implementation



Speeup: GPU vs CPU

Table: Overview of maximum achieved speed-ups – $\max(S_p(n_r))$ – for the two different GPU implementations, compared to the single-core CPU implementation.

| | naive | improved |
|------|-------|----------|
| GPU1 | 6.1× | 68× |
| GPU2 | 7.2× | 78× |
| GPU3 | 9× | 130× |